Relativistisk kvantemek 29.5.92

Løsninger

1a) Kanonisk form: Lineariserer $E^2 = p^2c^2 + (mc^2)^2$ og med minimal kobling $\vec{p} \to \vec{p} - e\vec{A}$

$H\psi = [c\vec{\alpha}(\vec{p} - e\vec{A}) + \beta mc^2 + e\varphi]\psi = E\psi$ med operatorene $\vec{p} = -i\hbar\nabla$, $E = i\hbar\frac{\partial}{\partial t}$

Dirac-matrisene oppfyller antikommutator-reglene

$\{\alpha^i, \alpha^k\}_+ = 2\delta^{ik}$, $\{\beta, \alpha^i\} = 0$, $\beta^2 = 1$ $\quad i,k = 1,2,3$ $\quad$ Spinor: $\Psi = \begin{pmatrix}\psi_1\\ \psi_2\\ \psi_3\\ \psi_4\end{pmatrix}$

$\underline{H\psi = [c\vec{\alpha}(-i\hbar\nabla - e\vec{A}) + \beta mc^2 + e\varphi]\Psi = i\hbar\frac{\partial}{\partial t}\Psi}$

Kovariant form.

Flytter over $E \to E - e\varphi$, multipliserer med β og setter $\beta\alpha^k = \gamma^k$, $\beta = \gamma^0$

$\underline{[\gamma^\mu(p_\mu - eA_\mu) - mc^2]\Psi = 0}$ med $\{\gamma^\mu, \gamma^\nu\} = 2g^{\mu\nu}$

b) Fri partikkel i ro $\vec{p} = 0$ $\quad p^0 = \frac{E}{c} = \frac{1}{c}(mc^2 + c^2\cancel{(p)^2}) = mc$, $\vec{A} = 0$ $\varphi = 0$

$(c\gamma^0 p^0 - mc^2)\Psi = 0$ $\quad cp^0 = i\hbar\frac{\partial}{\partial t}$

i standardrepresentasjonen: $\gamma^0 = \begin{pmatrix}1&&&\\&1&&\\&&-1&\\&&&-1\end{pmatrix}$, $\quad \left(i\hbar\frac{\partial}{\partial t} - mc^2\right)\Psi_a = 0$, $\quad \left(-i\hbar\frac{\partial}{\partial t} - mc^2\right)\Psi_b = 0$

med $\Psi_a = \begin{pmatrix}\psi_1\\ \psi_2\end{pmatrix} = u_a\left(\frac{1}{2\pi\hbar}\right)^{3/2} e^{-\frac{i}{\hbar}mc^2 t}$ $\quad u_a = \begin{pmatrix}u_1\\ u_2\end{pmatrix}$, $u_b = \begin{pmatrix}u_3\\ u_4\end{pmatrix}$ konstanter

$\Psi_b = \begin{pmatrix}\psi_3\\ \psi_4\end{pmatrix} = u_b\left(\frac{1}{2\pi\hbar}\right)^{3/2} e^{\frac{i}{\hbar}mc^2 t}$

4 løsninger:

$\underline{\Psi_0^{(i)} = u^{(i)}\left(\frac{1}{2\pi\hbar}\right)^{3/2} e^{-\frac{i}{\hbar}E_0^{(i)} t}}$, $\quad E_0^{(i)} = \begin{cases} mc^2 & \text{for } i = 1,2 \\ -mc^2 & i = 3,4\end{cases}$

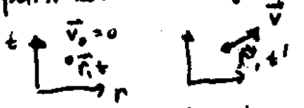
med $\underline{u_0^{(1)} = \begin{pmatrix}1\\0\\0\\0\end{pmatrix}, \quad u_0^{(2)} = \begin{pmatrix}0\\1\\0\\0\end{pmatrix}, \quad u_0^{(3)} = \begin{pmatrix}0\\0\\1\\0\end{pmatrix}, \quad u_0^{(4)} = \begin{pmatrix}0\\0\\0\\1\end{pmatrix}}$

$u^{(1)}$ og $u^{(3)}$ har spinnkomp. opp langs z-aksen, $u^{(2)}$ og $u^{(4)}$ ned:

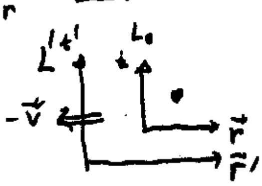
$\Sigma^z u_0^{(1)(3)} = \begin{pmatrix}\sigma^z & \\ & \sigma^z\end{pmatrix} u_0^{(1)(3)} = \begin{pmatrix}1&&&\\&-1&&\\&&1&\\&&&-1\end{pmatrix} u_0^{(1)(3)} = + u_0^{(1)(3)}$, $\quad \Sigma^z u_0^{(2)(4)} = + u_0^{(2)(4)}$

c) Sammenhengen $\Psi(\vec{r}', t') = S\,\Psi_0(\vec{r}, t)$ mellom tilstandsfunksjonen

for en partikkel i ro, $\Psi_0(\vec{r}, t)$, og en partikkel som går med hastighet $\vec{v}$, $\Psi(\vec{r}', t')$ er bestemt



av Lorentztransformasjonen mellom egensystemet L_0 og et inertialsystem som går med hastighet $-\vec{v}$.

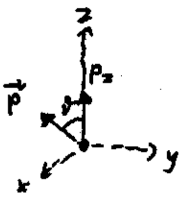


Ved en Lorentz-transformasjon og dreining er 4-skalar invariant $p_\mu x^\mu = E_0 t = p_\mu' x^{\mu'} = Et' - \vec{p}\vec{r}'$

Rel. kv.mek 29.5.92

Løsn., forts.

1c forts. Oppnår partikkel med $\vec{p}$ i retning $\vartheta, \varphi=0$ fra partikkel i ro ved en Lorentz-transf. langs z-aksen først og så dreining om y-aksen:



$$\psi(\vec{r}', t') = S^D S^L \psi_0(\vec{r}, t)$$

For partikkelhastighet langs z-aksen er

$$\gamma^0\gamma^3 = \beta\beta\alpha^3 = \alpha^3 = \begin{pmatrix} & \sigma^3 \\ \sigma^3 & \end{pmatrix} = \begin{pmatrix} & & 1 & \\ & & & -1 \\ 1 & & & \\ & -1 & & \end{pmatrix}$$

og

$$S^L = \begin{pmatrix} \cosh\frac{w}{2} & & \sinh\frac{w}{2} & \\ & \cosh\frac{w}{2} & & -\sinh\frac{w}{2} \\ \sinh\frac{w}{2} & & \cosh\frac{w}{2} & \\ & -\sinh\frac{w}{2} & & \cosh\frac{w}{2} \end{pmatrix} \qquad \cosh\frac{w}{2} = \sqrt{\frac{E+mc^2}{2mc^2}}, \quad \sinh\frac{w}{2} = \frac{E+mc^2}{2mc^2}\,\frac{c\,p_z}{E+mc^2}$$

Fra partikkel i ro med spinn opp langs z-aksen (tilstand $\psi_0^{(1)}$ før) får da tilstandsfunksjonen for partikkel med $\vec{p}=(0,0,p_z)$ og positiv helisitet.

$$\psi^L(\vec{r}_L, t_L) = S^L \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} \left(\frac{1}{2\pi\hbar}\right)^{3/2} e^{-\frac{i}{\hbar}E_0 t} = \begin{pmatrix} \cosh\frac{w}{2} \\ 0 \\ \sinh\frac{w}{2} \\ 0 \end{pmatrix} \left(\frac{1}{2\pi\hbar}\right)^{3/2} e^{-\frac{i}{\hbar}(E_L t_L - \vec{p}_L \vec{r}_L)}$$

$$= \sqrt{\frac{E+mc^2}{2mc^2}} \begin{pmatrix} 1 \\ 0 \\ \frac{c\,p_z}{E+mc^2} \\ 0 \end{pmatrix} \left(\frac{1}{2\pi\hbar}\right)^{3/2} e^{-\frac{i}{\hbar}(Et - p_z z)}$$

Videre til partikkel med $\vec{p} = (p_x, 0, p_z)$ i retning $\vartheta, \varphi = 0$.

$$\gamma^3\gamma^1 = \beta\alpha^3\beta\alpha^1 = -\alpha^3\alpha^1 = -\begin{pmatrix} & \sigma^3 \\ \sigma^3 & \end{pmatrix}\begin{pmatrix} & \sigma^1 \\ \sigma^1 & \end{pmatrix} = -\begin{pmatrix} \sigma^3\sigma^1 & \\ & \sigma^3\sigma^1 \end{pmatrix} = -i\begin{pmatrix} \sigma^2 & \\ & \sigma^2 \end{pmatrix} = -i\Sigma^2$$

$$= \begin{pmatrix} & \sigma^3 \\ -\sigma^3 & \end{pmatrix}\begin{pmatrix} & \sigma^1 \\ -\sigma^1 & \end{pmatrix} = \begin{pmatrix} & & 1 & \\ & & & -1 \\ -1 & & & \\ & 1 & & \end{pmatrix}\begin{pmatrix} & & & 1 \\ & & 1 & \\ & -1 & & \\ -1 & & & \end{pmatrix} = \begin{pmatrix} & -1 & & \\ 1 & & & \\ & & & -1 \\ & & 1 & \end{pmatrix}$$

$$S^D = \begin{pmatrix} \cos\vartheta/2 & -\sin\vartheta/2 & & \\ \sin\vartheta/2 & \cos\vartheta/2 & & \\ & & \cos\vartheta/2 & -\sin\vartheta/2 \\ & & \sin\vartheta/2 & \cos\vartheta/2 \end{pmatrix}$$

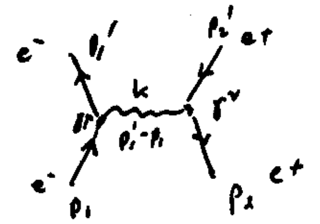
$$\psi(\vec{r}', t') = S^D \psi^L = \sqrt{\frac{E+mc^2}{2mc^2}} \left(\frac{1}{2\pi\hbar}\right)^{3/2} S^D \begin{pmatrix} 1 \\ 0 \\ \frac{c|\vec{p}|}{E+mc^2} \\ \end{pmatrix} e^{-\frac{i}{\hbar}(E_L t_L - p_L r_L)}$$

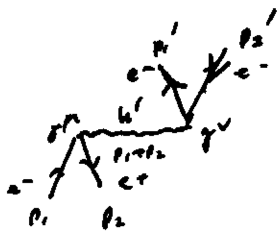
$$= \sqrt{\frac{E+mc^2}{2mc^2}} \begin{pmatrix} \cos\vartheta/2 \\ \sin\vartheta/2 \\ \frac{c|\vec{p}|}{E+mc^2}\cos\vartheta/2 \\ \frac{c|\vec{p}|}{E+mc^2}\sin\vartheta/2 \end{pmatrix} \left(\frac{1}{2\pi\hbar}\right)^{3/2} e^{-\frac{i}{\hbar}(Et - \vec{p}\vec{r})} \qquad \vec{p} = (p_x, 0, p_z)$$

Rel. kv.mek 29.5.97

Løsn. forts

1 d)





i siste diagram har p_1' og p_2 byttet plass så et − tegn.

Bruker $\hbar = c = 1$

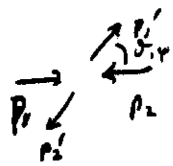
$$S = \left(\frac{1}{2\pi}\right)^{3/2\cdot 4} \frac{m^{1/2\cdot 4}}{(E_1'E_1E_2'E_2)^{1/2}} (-ie^2)^2 (2\pi)^{4\cdot 2} \frac{-i}{(2\pi)^4} \int d^4k \, (\bar{u}_1' \gamma^\mu u_1) \frac{g_{\mu\nu}}{k^2} (\bar{v}_2 \gamma^\nu v_2') \delta(p_1-p_1'-k)\,\delta(p_2-k-p_2')$$

$$- \left(\frac{1}{2\pi}\right)^6 \frac{m^2}{(E_1'E_1E_2'E_2)^{1/2}} (-e^2)(2\pi)^8 \frac{-i}{(2\pi)^4} \int d^4k' \, (\bar{u}_1' \gamma^\nu v_2') \frac{g_{\mu\nu}}{k'^2} (\bar{v}_2 \gamma^\mu u_1) \delta(p_1+p_2-k')\,\delta(p_1'+p_2'-k')$$

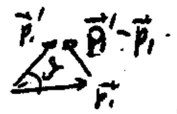
$$= \left(\frac{1}{2\pi}\right)^2 \frac{m^2\, ie^2}{(E_1'E_1E_2'E_2)^{1/2}} \left[\frac{(\bar{u}_1'\gamma^\mu u_1)(\bar{v}_2\gamma_\mu v_2')}{(p_1-p_1')^2} - \frac{(\bar{u}_1'\gamma^\mu v_2')(\bar{v}_2\gamma_\mu u_1)}{(p_1+p_2)^2} \right] \delta(p_1'+p_2'-p_1-p_2)$$

1e I tyngdepunktsystemet $\vec{p}_2 = -\vec{p}_1$, $E_1 = E_2 = E$ Setter

$\vec{p}_2\,' = -\vec{p}_1\,'$, $E_1' = E_2' = E$



Regner ut $\dfrac{(\bar{u}_1'\gamma^1 u_1)(\bar{v}_2\gamma_1 v_2')}{(p_1-p_1')^2}$



$$(p_1-p_1')^2 = (E_1-E_1')^2 - (\vec{p}_1-\vec{p}_1\,')^2 = -\left(2|\vec{p}|\sin\vartheta/2\right)^2$$

e^- inn: $\vartheta = 0$ (med $\hbar = c = 1$)
pos. helisitet

$$u_1 = \sqrt{\frac{E+m}{2m}} \begin{pmatrix} 1 \\ 0 \\ \frac{|\vec{p}|}{E+m} \\ 0 \end{pmatrix} e^{-i\varphi/2}$$

e^- ut: ϑ, φ

$$u_1' = \sqrt{\frac{E+m}{2m}} \begin{pmatrix} \cos\vartheta/2 \; e^{-i\varphi/2} \\ \sin\vartheta/2 \; e^{i\varphi/2} \\ \frac{|\vec{p}|}{E+m} \cos\vartheta/2 \; e^{-i\varphi/2} \\ \frac{|\vec{p}|}{E+m} \sin\vartheta/2 \; e^{i\varphi/2} \end{pmatrix}$$

e^+ inn: $\vartheta = \pi$, $\varphi+\pi$

$$v_2 = \sqrt{\frac{E+m}{2m}} \begin{pmatrix} -\frac{|\vec{p}|}{E+m} \\ 0 \\ 1 \\ 0 \end{pmatrix} (-i) e^{-i\varphi/2}$$

e^+ ut: $\pi-\vartheta$, $\varphi+\pi$

$$v_2' = \sqrt{\frac{E+m}{2m}} \begin{pmatrix} \frac{|\vec{p}|}{E+m} \cos\vartheta/2 \; i e^{-i\varphi/2} \\ \frac{|\vec{p}|}{E+m} \sin\vartheta/2 \; i e^{i\varphi/2} \\ -\cos\vartheta/2 \; i e^{-i\varphi/2} \\ -\sin\vartheta/2 \; i e^{i\varphi/2} \end{pmatrix}$$

$$\gamma^0\gamma^1 = \begin{pmatrix} 1 & & & \\ & 1 & & \\ & & -1 & \\ & & & -1 \end{pmatrix} \begin{pmatrix} & & & 1 \\ & & 1 & \\ & -1 & & \\ -1 & & & \end{pmatrix} = \begin{pmatrix} & & & 1 \\ & & 1 & \\ & 1 & & \\ 1 & & & \end{pmatrix} \qquad \gamma^0\gamma_1 = -\begin{pmatrix} & & & 1 \\ & & 1 & \\ & 1 & & \\ 1 & & & \end{pmatrix}$$

$$\bar{u}_1'\gamma^1 u_1 = u_1'^{+}\gamma^0\gamma^1 u_1 = u_1'^{+}\begin{pmatrix} & & & 1 \\ & & 1 & \\ & 1 & & \\ 1 & & & \end{pmatrix} u_1 = \frac{E+m}{2m}\left[\frac{|\vec{p}|}{E+m}\sin\vartheta/2 \; e^{-i\varphi/2} e^{-i\varphi/2} + \sin\vartheta/2 \; e^{-i\varphi/2} \frac{|\vec{p}|}{E+m} e^{-i\varphi/2}\right]$$

$$= \underline{\frac{|\vec{p}|}{m}\sin\vartheta/2 \; e^{-i\varphi}}$$

$$\bar{v}_2\gamma_1 v_2' = -v_2^{+}\begin{pmatrix} & & & 1 \\ & & 1 & \\ & 1 & & \\ 1 & & & \end{pmatrix} v_2' = -\frac{E+m}{2m}\left[\frac{-|\vec{p}|}{E+m}\, i e^{+i\varphi/2} (-i)\sin\vartheta/2 \; e^{i\varphi/2} + (i) e^{i\varphi/2} \frac{|\vec{p}|}{E+m}\sin\vartheta/2 \; i e^{i\varphi/2}\right]$$

$$= \frac{|\vec{p}|}{m}\sin\vartheta/2 \; e^{i\varphi}$$

$$\frac{ie^2}{(2\pi)^2}\left(\frac{m}{E}\right)^2 \frac{\left(\frac{|\vec{p}|}{m}\right)^2 \sin^2\vartheta/2 \; e^{-i\varphi} e^{i\varphi}}{-4|\vec{p}|^2\sin^2\vartheta/2} = \underline{\underline{-\frac{ie^2}{(2\pi)^2\, 4E^2}}}$$

Rel. kv. mek. 29.3.92

Em. felt

2a) $\mathcal{L} = -\frac{1}{2\mu_0} K^{\mu\nu\rho\sigma}(\partial_\mu A_\nu)(\partial_\rho A_\sigma) = -\frac{1}{2\mu_0}\left[(\partial^\mu A^\nu)(\partial_\mu A_\nu) - (\partial^\nu A^\mu)(\partial_\mu A_\nu)\right] = -\frac{1}{2\mu_0} F^{\mu\nu}\partial_\mu A_\nu$

$= -\frac{1}{4\mu_0}\left(F^{\mu\nu}\partial_\mu A_\nu + F^{\nu\mu}\partial_\nu A_\mu\right) = -\frac{1}{4\mu_0} F^{\mu\nu}F_{\mu\nu} = \frac{1}{2\mu_0}(\vec{E}\vec{D} - \vec{B}\vec{H})$

Feltlikn:

$0 = \frac{\partial \mathcal{L}}{\partial A_\lambda} - \partial_\kappa \frac{\partial \mathcal{L}}{\partial(\partial_\kappa A_\lambda)} = 0 - \partial_\kappa\left[-\frac{1}{2\mu_0}(g^{\mu\rho}g^{\nu\sigma} - g^{\nu\rho}g^{\mu\sigma})\left[\delta_{\mu\kappa}\delta_{\nu\lambda}(\partial_\rho A_\sigma) + (\partial_\mu A_\nu)\delta_{\rho\kappa}\delta_{\sigma\lambda}\right]\right]$

$= \frac{1}{2\mu_0}\partial_\kappa\left[(\partial^\kappa A^\lambda - \partial^\lambda A^\kappa) + (\partial^\kappa A^\lambda - \partial^\lambda A^\kappa)\right] = \underline{\frac{1}{\mu_0}\partial_\kappa F^{\kappa\lambda} = 0}$

De homogene Maxwells likninger følger fra formen til $F^{\mu\nu} = \partial^\mu A^\nu - \partial^\nu A^\mu$

$$\Rightarrow \quad \underline{\varepsilon^{\kappa\lambda\mu\nu}\partial_\lambda F_{\mu\nu} = 0}$$

Gauge-invarians:

$F^{\mu\nu} = \partial^\mu A^\nu - \partial^\nu A^\mu$ forandrer ikke verdi med ny $A^{\mu\prime} = A^\mu + \partial^\mu \chi$

hvor $\chi(\vec{r},t)$ er vilkårlig funksjon!

$\underline{F^{\mu\nu\prime}} = \partial^\mu(A^\nu + \partial^\nu\chi) - \partial^\nu(A^\mu + \partial^\mu\chi) = \partial^\mu A^\nu - \partial^\nu A^\mu + \partial^\mu\partial^\nu\chi - \partial^\nu\partial^\mu\chi = \underline{F^{\mu\nu}}$

Hvis Lorentz bet. ikke oppfylt: $\partial_\mu A^\mu = B \neq 0$ så velger ny $A^{\mu\prime} = A^\mu + \partial^\mu\chi$

med χ som løsning av $\partial_\mu\partial^\mu\chi = -B$ som har løsning

Da blir $\underline{\partial_\mu A^{\mu\prime}} = \partial_\mu A^\mu + \partial_\mu\partial^\mu\chi = B - B = \underline{0}$ med samme $F^{\mu\nu}$

b) Kanoniske impulser:

$\pi^\kappa = \frac{\partial\mathcal{L}}{\partial \dot{A}_\kappa} = \frac{1}{c}\frac{\partial\mathcal{L}}{\partial(\partial_0 A_\kappa)} = -\frac{1}{\mu_0 c}(\partial^0 A^\kappa - \partial^\kappa A^0) = \underline{\frac{1}{\mu_0 c}F^{\kappa 0}}$ $\qquad \pi^0 = \frac{1}{\mu_0 c}F^{00} = 0$

$\pi^k = \frac{1}{\mu_0 c}F^{k0} = \frac{1}{\mu_0}E^k = D^k$

$\mathcal{H} = \pi^\mu \dot{A}_\mu - \mathcal{L} = -\frac{1}{\mu_0 c}(\partial^0 A^\mu - \partial^\mu A^0)\, c\,\partial_0 A_\mu + \frac{1}{2\mu_0}(\partial_\nu A^\mu \partial^\nu A_\mu - \partial_\nu A^\mu \partial_\mu A^\nu)$

$= -\frac{1}{\mu_0}(\partial^0 A^\mu)(\partial_0 A_\mu) + \frac{1}{2\mu_0}(\partial^\nu A^\mu)(\partial_\nu A_\mu) + \frac{1}{\mu_0}(\partial^\mu A^0)(\partial_0 A_\mu) - \frac{1}{2\mu_0}(\partial^\nu A^\mu)(\partial_\mu A_\nu)$

De to første leddene vil gi den oppgitte integranden i H

De to siste former om til totale differential ved å benytte Lorentz bet. $\partial^\nu A_\nu = 0$: $(\partial^\nu A^\mu)(\partial_\mu A_\nu) + A^\mu \underbrace{\partial_\mu\partial^\nu A_\nu}_{=0} = \partial^\nu(A^\mu \partial_\mu A_\nu)$

$\mathcal{H} = -\frac{1}{2\mu_0}\partial^0 A^\mu \partial_0 A_\mu + \frac{1}{2\mu_0}\partial^k A^\mu \partial_k A_\mu + \frac{1}{\mu_0}\partial^\mu(A^0\partial_0 A_\mu) - \frac{1}{2\mu_0}\partial^\nu(A^\mu\partial_\mu A_\nu)$

som gir ($\partial_k = -\partial^k$ $\quad \partial_0 = \partial^0 = \frac{1}{c}\frac{\partial}{\partial t}$)

$H = \int d^3x\, \mathcal{H} = \underline{-\frac{1}{2\mu_0}\int d^3x\left[\frac{1}{c^2}\dot{A}^\mu \dot{A}_\mu + \partial_k A^\mu \partial_k A_\mu\right]}$ $\quad$ da $\int d^3x\, \partial^\mu(A^\mu\partial_\mu A_\nu) = (A^\mu\partial_\mu A)_{\text{overflate}} = 0$

· Lorentz betingelsen anvendt på plan bølge oppløsningen gir

$\partial_\mu A^\mu = \int \frac{d^3k}{(2\pi)^{3/2}}\left(\frac{\hbar}{2\varepsilon_0\omega_k}\right)^{1/2}\left[-ik_\mu a^\mu(k)e^{-ik_\nu x^\nu} + ik_\mu a^{\mu\dagger}(k)e^{ik_\nu x^\nu}\right] = 0 \Rightarrow \underline{k_\mu a^\mu = 0}$

For bølge i retning $\vec{k}$ (f.eks i z-retningen $\vec{k} = (0,0,k^3)$)

$k_\mu a^\mu = k^0 a^0 - \vec{k}\vec{a} = \frac{\omega_k}{c}a^0 - k\,a^\parallel = 0$ $\quad$ mens $\vec{a}_\perp$ er fri

(Med $\vec{k} = 0,0,k^3$) $\quad \frac{\omega_k}{c}a^0 - k^3 a^3 = 0$ $\qquad a^1$ og a^2 er fri)

$\frac{\omega_k}{c} = |\vec{k}| = |k^3|$ $\qquad a^0 \mp a^3 = 0$ $\qquad$ Feltets longitudinelle og tidskomponenter bestemmes hverandre.

Rel. kv. mek. 29.5.92
Løsn. fortsatt

2c) $$[A_\mu(x), \varepsilon_0 \dot{A}_\nu(y)]_{x_0=y_0} = \int \frac{d^3k\, d^3k'}{(2\pi)^3} \frac{\varepsilon_0 \hbar (-i\omega_{k'})}{2\varepsilon_0 (\omega_k \omega_{k'})^{1/2}} \Big[[a_\mu(k), a_\nu(k')] e^{-i(kx+k'y)} + [a_\mu(k), a_\nu^+(k')] e^{-i(kx-k'y)} + [a_\mu^+(k), a_\nu(k')] e^{i(kx-k'y)} - [a_\mu^+(k), a_\nu^+(k')] e^{i(kx+k'y)} \Big]_{x_0=y_0}$$

$$= \int \frac{d^3k\, d^3k'}{(2\pi)^3} \frac{-i\hbar}{2} \left(\frac{\omega_{k'}}{\omega_k}\right) \left[+g_{\mu\nu}\, \delta(k-k') e^{-i(kx+k'y)} + g_{\mu\nu}\, \delta(k-k') e^{i(kx-k'y)} \right]_{x_0=y_0}$$

$$= -\frac{i\hbar}{2} g_{\mu\nu} \left(\frac{1}{2\pi}\right)^3 \int d^3k \left[e^{-i\vec{k}(\vec{x}-\vec{y})} + e^{i\vec{k}(\vec{x}-\vec{y})} \right] \quad \text{når } x_0 = y_0$$

$$= \underline{-i\hbar\, g_{\mu\nu}\, \delta(\vec{x}-\vec{y})}$$

2d) For det kvantiserte feltet gir Lorentz-betingelsen, som generell likning for feltoperatorene $A^\mu(x)$, inkonsistens:
Ser dette ved å anvende $\frac{\partial}{\partial x_\mu}$ på kommutatoren ovenfor
Venstre side:
$$\frac{\partial}{\partial x_\mu} [A_\mu(x), \varepsilon_0 \dot{A}_\nu(y)] = [\partial^\mu A_\mu(x), \varepsilon_0 \dot{A}_\nu(y)] = 0 \quad \text{hvis } \partial_\mu A^\mu(x) = 0 \text{ generelt}$$

Høyre side:
$$-i\hbar\, g_{\mu\nu} \frac{\partial}{\partial x_\mu} \delta(\vec{x}-\vec{y}) \neq 0 \qquad \left(\int_{-\varepsilon}^{+\varepsilon} f(x)\, \delta'(x)\, dx = -f(0) \neq 0 \text{ generelt.} \right)$$

Kan ordne dette ved forandre kommutator-reglen til bare å gjelde feltet "transversale komponenter".
Annen måte ved å innskrenke mulige felttilstander:
Tillater bare tilstander $|\psi\rangle$ som oppfyller $\underline{\partial_\mu A^{\mu(+)} |\psi\rangle = 0}$
Her er $A^{\mu(+)}$ den del av A^μ som inneholder annihilasjonsoperatorer.
Herav følger $\langle\psi| \partial_\mu A^{\mu(-)} = 0$ og dermed for alle tillatte tilstander
$$\langle\psi| \partial_\mu A^\mu |\psi\rangle = \langle\psi| \partial_\mu A^{\mu(-)}(x) + \partial_\mu A^{\mu(+)}(x) |\psi\rangle = \langle\psi|\partial_\mu A^{\mu(-)}|\psi\rangle + \langle\psi|\partial_\mu A^{\mu(+)}|\psi\rangle = 0$$

Fra spektral-oppløsningen følger da
$$\partial_\mu A^{\mu(+)}|\psi\rangle = \int \frac{d^3k}{(2\pi)^{3/2}} \left(\frac{\hbar}{2\varepsilon_0 \omega_k}\right)^{1/2} (-ik_\mu a^\mu)\, e^{-ikx} |\psi\rangle = 0 \quad \Rightarrow \quad \underline{k_\mu a^\mu(k)|\psi\rangle = 0}$$

For $k = (\frac{\omega_k}{c}, \vec{k}) = (\frac{\omega_k}{c}, 0, 0, k^3) \Rightarrow (\frac{\omega_k}{c} a^0 - \vec{k}\vec{a})|\psi\rangle = (\frac{\omega_k}{c} a^0 - k^3 a^3)|\psi\rangle = 0$

$(\frac{\omega_k}{c})^2 = (\vec{k})^2 \Rightarrow \omega_k = ck^3 \Rightarrow \underline{(a^0 - a^3)|\psi\rangle = 0}$

Energi: For hver frekvenskomponent får vi (når vi bruker normalordning)
$$\langle\psi|H|\psi\rangle = \frac{1}{2} \int d^3k\, \hbar\omega_k :(a_\mu(k)\, a^{\mu+}(k) + a_\mu^+(k)\, a^\mu(k)): = -\int d^3k\, \hbar\omega_k \langle\psi| a_\mu^+(k)\, a^\mu(k)|\psi\rangle$$

$$\langle\psi|a_\mu^+(k)\, a^\mu(k)|\psi\rangle = \langle\psi| a^{0+}a^0 - a^{1+}a^1 - a^{2+}a^2 - a^{3+}a^3|\psi\rangle = -\langle\psi| a^{1+}a^1 + a^{2+}a^2|\psi\rangle$$

idet
$$\langle\psi| a^{0+}a^0 - a^{3+}a^3|\psi\rangle = \frac{1}{2}\left(\langle\psi| (a^{0+} + a^{3+}) \underbrace{(a^0 - a^3)|\psi\rangle}_{=0} + \underbrace{\langle\psi|(a^{0+} - a^{3+})}_{=0} (a^0 + a^3)|\psi\rangle \right) = \underline{0}$$

Altså:
$$\underline{E = \langle\psi|H|\psi\rangle = \int d^3k\, \hbar\omega_k \langle\psi|a_1^+ a_1 + a_2^+ a_2|\psi\rangle \geq 0}$$ for alle tilstandene som oppfyller den modifiserte Lorentz-betingelsen $\underline{\partial_\mu A^{\mu(+)}|\psi\rangle = 0}$