

1) E 2) $M_{JT} = 0$ 3) Syklisk regel $\Leftrightarrow \frac{\alpha_p K_T}{\alpha_V} - \frac{1}{P} \Leftrightarrow \underline{\underline{P = \frac{RT}{V}}}$

4) $K_T = -\frac{1}{V} \left(\frac{\partial V}{\partial P} \right)_{T,N}$

$$\frac{P}{kT} = \frac{N}{V} + \frac{N^2}{2V^2}$$

$$\frac{dP}{kT} - \frac{P dT}{kT^2} = \frac{dN}{V} - \frac{N dV}{V^2} + \frac{N dN}{V^2} - \frac{N^2 dV}{V^3}$$

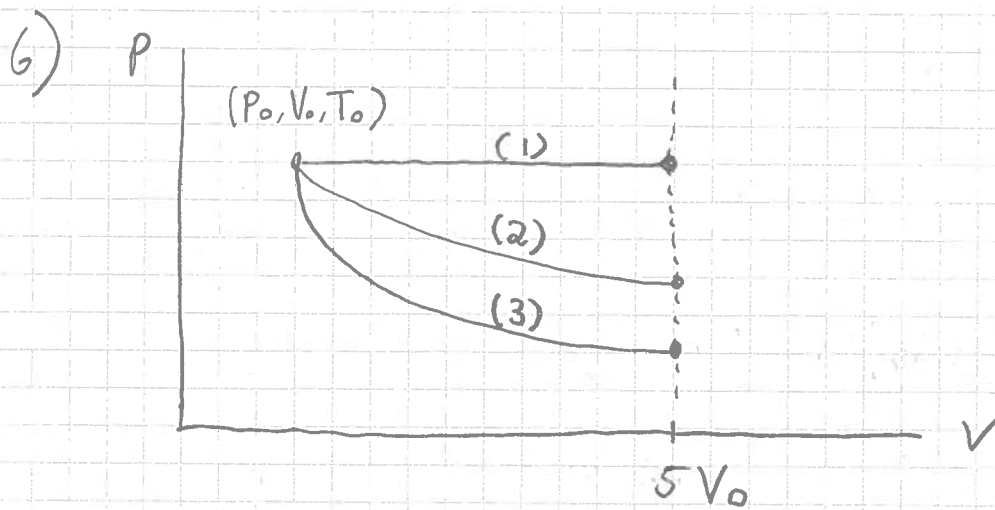
konstant T og N:

$$\frac{dP}{kT} = - \frac{N dV}{V^2} \left(1 + \frac{N}{V} \right)$$

$$\Rightarrow \left(\frac{\partial V}{\partial P} \right)_{T,N} = \frac{-1}{kT} \frac{1}{\frac{N}{V^2} \left(1 + \frac{N}{V} \right)}$$

$$\Rightarrow \underline{\underline{K_T = \frac{1}{kT} \frac{V}{N} \left(\frac{1}{1 + N/V} \right)}}$$

5) Arbejd = all energi som bruges gennetlæret Sam
ikke skyldes ΔT .



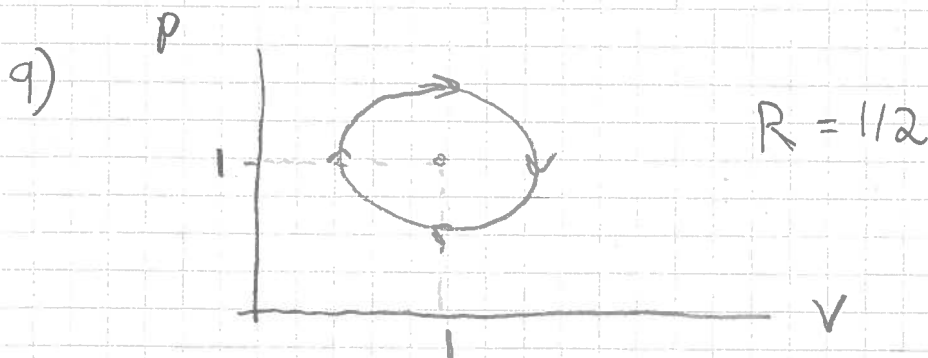
$$\underline{W_P > W_T > W_S}$$

7) $C_P - C_V = Nk$

$$\underline{C_P = \frac{5}{2} Nk}$$

8) Adiabatic = constant entropy

$$\Rightarrow \underline{\Delta S = 0}$$



$$W = \pi R^2 = \underline{\underline{\frac{\pi}{4}}}$$

$$10) \eta_c = 1 - \frac{T_{\text{minimum}}}{T_{\text{maksimum}}}$$

I deell gass $PV = nRT$

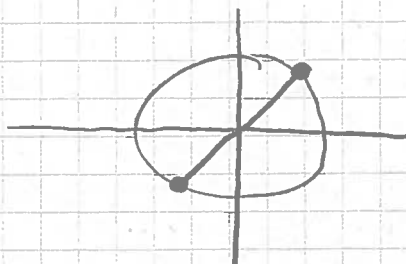
$$nRT = \left(1 + \frac{1}{2} \sin \varphi\right) \left(1 + \frac{1}{2} \cos \varphi\right)$$

För å finne $T_{\text{max/min}}$ må vi løse $dT/d\varphi = 0$

$$\begin{aligned} nR \frac{dT}{d\varphi} &= \left(\frac{1}{2} \cos \varphi\right) \left(1 + \frac{1}{2} \cos \varphi\right) + \left(1 + \frac{1}{2} \sin \varphi\right) \left(-\frac{1}{2} \sin \varphi\right) \\ &= \frac{1}{2} (\cos \varphi - \sin \varphi) + \frac{1}{4} (\cos^2 \varphi - \sin^2 \varphi) = 0 \end{aligned}$$

$$T_{\text{max}} : \varphi = \pi/4$$

$$T_{\text{min}} : \varphi = 5\pi/4$$



$$nRT_{\text{min}} = \left(1 - \frac{1}{2\sqrt{2}}\right)^2$$

$$nRT_{\text{max}} = \left(1 + \frac{1}{2\sqrt{2}}\right)^2$$

$$\eta_c = 1 - \frac{T_{\text{min}}}{T_{\text{max}}} \approx \underline{\underline{77\%}}$$

$$11) \Delta S = \text{tilst. funksjon} = 0$$

$$12) dS > 0 \text{ i irrev. prosess.}$$

13) En isotrop funktion er retningssuavhengig

$$F(x, y, z) = F(\vec{r}) = F(|\vec{r}|)$$

$$F(v_x, v_y, v_z) = A(v_x^2 + v_y^2 + v_z^2)$$

$$14) v_{rms} = \sqrt{\langle v^2 \rangle} = \left(\int_0^\infty v^2 f(v) dv \right)^{1/2}$$

$$= \dots = \underline{\underline{\sqrt{\frac{3kT}{m}}}}$$

15) C

$$16) 2D : f_{translasjon} = 2$$

$$f_{rot} = 1$$

$$\left. \begin{array}{l} f_{vib} = 1 \\ f_{pot. energi} = 1 \end{array} \right\} \begin{array}{l} \text{bidrar ikke ved} \\ \text{romtemperatur} \end{array}$$

$$C_v = \frac{1}{2} k (2 + 1) = \underline{\underline{\frac{3}{2} k}}$$

$$17) F = NkT \ln\left(\frac{N\lambda^3}{V}\right) - NkT$$

$$= -NkT \ln V + \text{irrelevante led} \\ (\text{uten } V)$$

$$P = -\left(\frac{\partial F}{\partial V}\right)_T = \underline{\underline{\frac{NkT}{V}}} = \text{ideell gass}$$

$$18) P\sqrt{V} = AT$$

$$H = U + PV = PV \Rightarrow U = 0 \Rightarrow \underline{C_v = 0}$$

$$C_p - C_v = T \left(\frac{\partial P}{\partial T}\right)_V \left(\frac{\partial V}{\partial T}\right)_P$$

$$= T \frac{\partial}{\partial T} \left\{ \frac{AT}{\sqrt{V}} \right\} \frac{\partial}{\partial T} \left\{ \frac{A^2 T^2}{P^2} \right\}$$

$$= T \left\{ \frac{A}{\sqrt{V}} \right\} \left\{ 2 \frac{A^2 T}{P^2} \right\} = \left\{ \frac{A}{\sqrt{V}} \right\} 2 \{ V \}$$

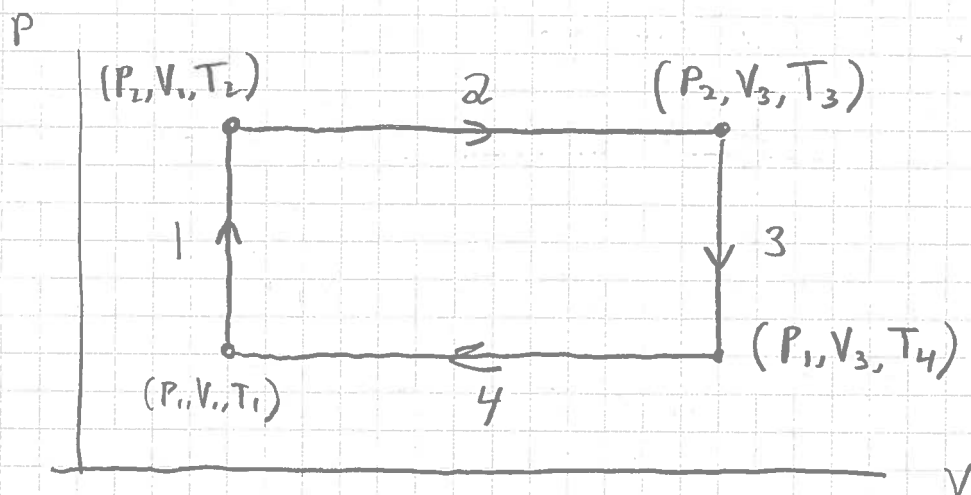
$$= \underline{\underline{2A\sqrt{V}}}$$

integrert over alternativ rev. prosess.

$$19) \Delta S = \int dS = \int_0^{V_1} \underbrace{C_v}_{A\sqrt{V}} \frac{dT}{T} + \underbrace{\left(\frac{\partial P}{\partial T}\right)_V}_{\frac{A}{\sqrt{V}}} dV$$

$$= 2A \left[\sqrt{V} \right]_{V_0}^{V_1} = \underline{\underline{2A(\sqrt{V_1} - \sqrt{V_0})}}$$

20)



$$dQ = dU + p dV$$

$$\text{Isochor: } dQ = dU \Rightarrow \Delta Q = C_v \Delta T$$

$$\text{Isobar: } dQ = dU + \underbrace{p dV}_{\text{konstant}} \Rightarrow \Delta Q = C_v \Delta T + p \Delta V$$

$$\Delta Q = C_v \Delta T + nR \Delta T = \Delta T \underbrace{(C_v + nR)}_{C_p}$$

$$= C_p \Delta T$$

$$Q_1 = C_v (T_2 - T_1) > 0$$

$$Q_2 = C_p (T_3 - T_2) > 0$$

$$Q_3 < 0$$

$$Q_4 < 0$$

$$W = \text{Areal} = (P_2 - P_1)(V_3 - V_1)$$

$$\eta = \frac{W}{Q_{\text{inn}}} = \frac{(P_2 - P_1)(V_3 - V_1)}{C_V(T_2 - T_1) + C_P(T_3 - T_2)}$$

21) $dS > 0 \Leftrightarrow$ minimier F für ä. offene LV

$\Rightarrow$ Tilst. lign. følger fra $dF = 0$

$$A \Leftrightarrow B \Leftrightarrow E$$

$$22) \quad dS = \underbrace{C_V \frac{dT}{T}}_{\text{Isoterm} = 0} + \underbrace{\left(\frac{\partial P}{\partial T}\right)_V}_{\frac{nR}{V}} dV$$

$$\Delta S = \int dS = nR \int_{V_1}^{V_2} \frac{dV}{V} = \underline{\underline{nR \ln \frac{V_2}{V_1} > 0}}$$

30. juni
2021

23) Siden vannmengden har samme temperatur T_0 er det ingen temperatur-utjevning i prosessen.

Da prosessen er irreversibel må vi integre langs en alternativ rev. vei (Isobar)

$$\Delta S_1 = \int_{T_0}^{T_0} \frac{C_p m dT}{T} = 0$$

$$\Delta S_2 = \int_{T_0}^{T_0} \frac{C_p m dT}{T} = 0$$

Adiabatisk:

$$\Delta S_{\text{univers}} = \Delta S_1 + \Delta S_2 = \underline{\underline{0}}$$

Det blir kun ~~en~~ ΔS hvis
temperaturene er forskjellige.

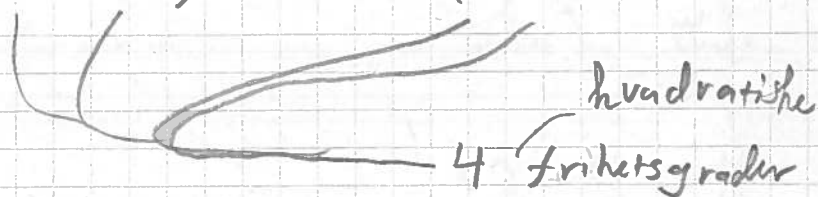
24) Ingen

$$25) \quad dF = -S dT - P dV$$

$$S = -\left(\frac{\partial F}{\partial T}\right)_V \quad \& \quad P = -\left(\frac{\partial F}{\partial V}\right)_T$$

$$\underline{\underline{\frac{\partial^2 F}{\partial V \partial T} = \left(\frac{\partial S}{\partial V}\right)_T = \left(\frac{\partial P}{\partial T}\right)_V}}$$

$$26) \quad E = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2) + \frac{1}{2} m \omega^2 (x^2 + y^2)$$


 kvadratiske
4 frihetsgrader

$$C_V = \frac{4}{2} k = \underline{\underline{2k}}$$

$$27) \quad \underline{\underline{\Delta S = 0}}$$

Lige mængde tilstande med

$$E = -J/2$$

Sam med $E = +J/2$

$$28) \quad C_H = \left(\frac{\partial H}{\partial T}\right)_H \quad \text{entropi}$$

$$H = U - \mu_0 H M = -\mu_0 H M$$

$$= -\mu_0 A \frac{H^2}{T}$$

$$\underline{\underline{C_H = \mu_0 A \left(\frac{H}{T}\right)^2}}$$

29) $\subset$ Dette er magnetisk binding.

30) Når $T = \infty$ er den termiske energien ∞ .

Da er det like sannsynlig å være i tilstande

~~$E = E_1, E_2, E_3$~~ E_1, E_2 , og E_3 .

$$P_3 = P_2 = P_1 = \underline{\underline{1/3}}$$

31) $dw = p dv$

$$H = U + PV$$

$$dH = \underbrace{du + p dv}_{T ds} + v dp$$

Naturlige variable er S og P .

Med $dw = y dx$:

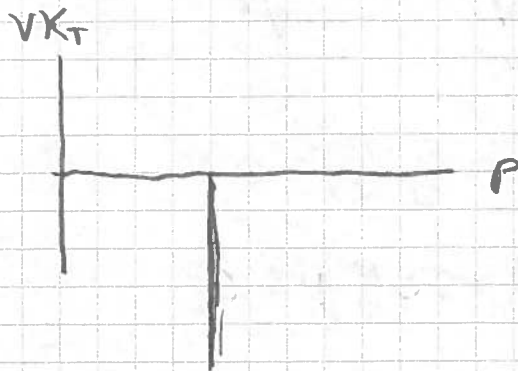
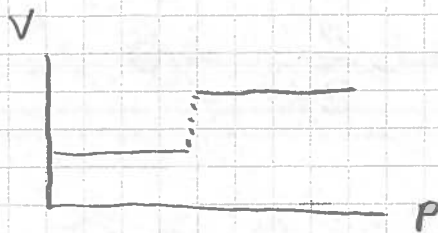
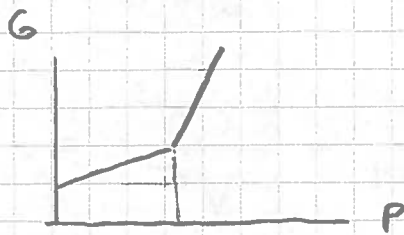
————— || ————— S og y

32) $f = c + 2 - q$

$$= 3 + 2 - 3 = \underline{\underline{2}}$$

$$33) \quad \kappa_T V = - \left(\frac{\partial V}{\partial P} \right)_T$$

$$V = \left(\frac{\partial G}{\partial P} \right)_T$$



A

$$34) P = \frac{kTN}{V} - \frac{b}{2} \frac{N^2}{V^2} + \frac{cN^3}{6V^3}$$

$$\frac{\partial P}{\partial V} = -\frac{kTN}{V^2} + b \frac{N^2}{V^3} - \frac{cN^3}{2V^4} = 0 \quad (1)$$

$$\frac{\partial^2 P}{\partial V^2} = 2 \frac{kTN}{V^3} - 3b \frac{N^2}{V^4} + \frac{2cN^3}{V^5} = 0 \quad (2)$$

$$(1): -kT + b \frac{N}{V} - c \frac{N^2}{2V^2} = 0$$

$$(2): 2kT - 3b \frac{N}{V} + 2c \frac{N^2}{V^2} = 0$$

$$(1): -kT + \frac{N}{V} \left(b - \frac{c}{2} \frac{N}{V} \right) = 0$$

$$(2): 2kT + \frac{N}{V} \left(-3b + 2c \frac{N}{V} \right) = 0$$

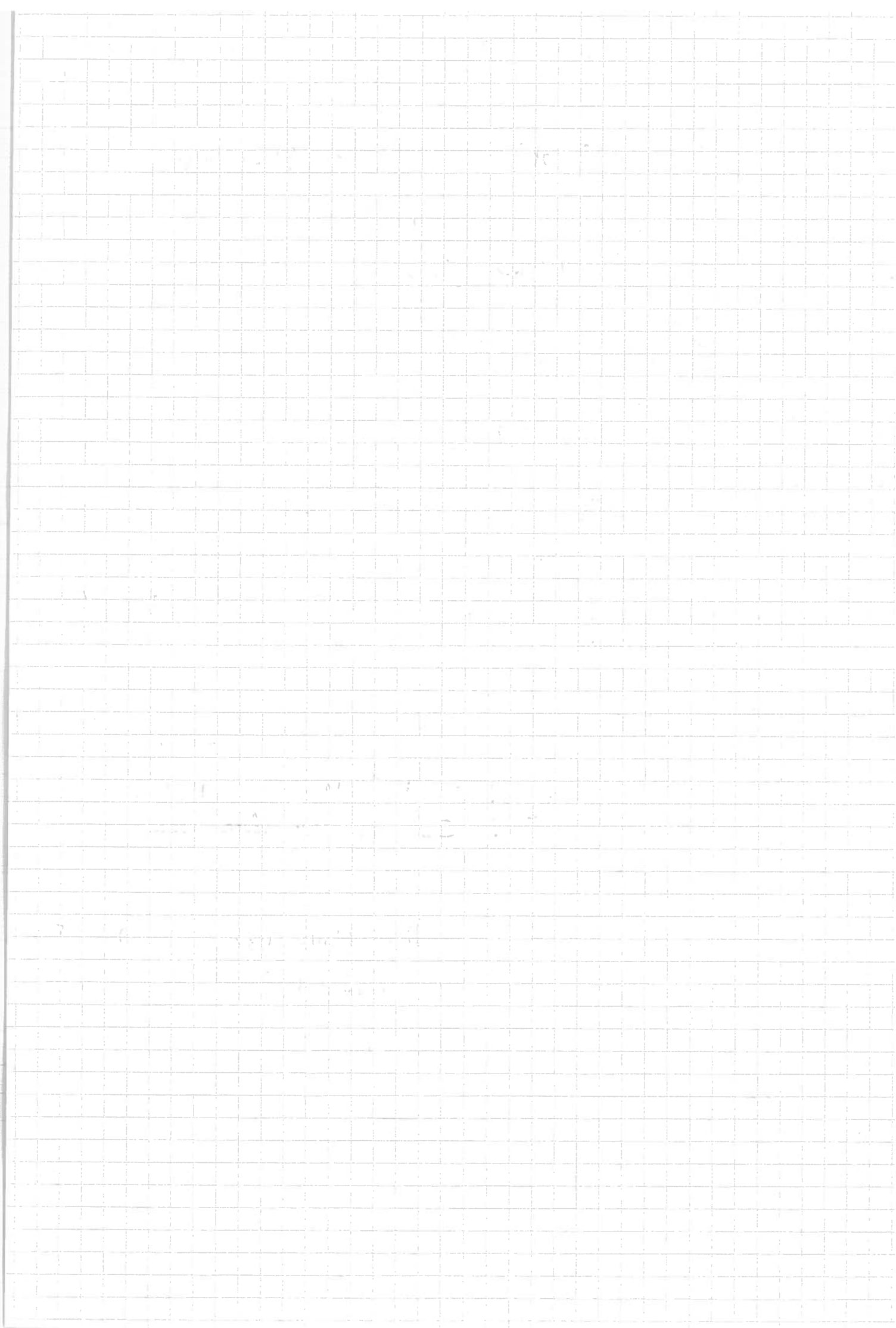
34) $P = kTn - \frac{b}{2}n^2 + \frac{c}{6}n^3$

$$\frac{\partial P}{\partial n} = kT - bn + \frac{c}{2}n^2 = 0 \quad (a)$$

$$\frac{\partial^2 P}{\partial n^2} = -b + cn = 0 \rightarrow b = cn \quad (b)$$

$$(a): kT - cn^2 + \frac{c}{2}n^2 = 0$$

$$kT = \frac{c}{2}n^2$$



$$34) \quad p = \frac{NkT}{V-Nb} - a \frac{N^2}{V^2}$$

b : Volumet okkupert av et gassmolekyl

a : et mål på den attraktive vekselvirkningen mellom gassmolekylene.

B

$$35) \quad p = \frac{NkT}{V} - \frac{aN^2}{V^2}$$

$$\frac{\partial p}{\partial V} = -\frac{NkT}{V^2} + 2 \frac{aN^2}{V^3} = 0 \Rightarrow -kT + 2 \frac{aN}{V} = 0 \quad (a)$$

$$\frac{\partial^2 p}{\partial V^2} = 2 \frac{NkT}{V^3} - 6 \frac{aN^2}{V^4} = 0 \Rightarrow 2kT - 3 \frac{aN}{V} = 0 \quad (b)$$

$$(a): \quad a = \frac{kT_c V_c}{2N}$$

$$(b): \quad a = \frac{kT_c V_c}{3N}$$

} eneste mulighet å oppfylle begge
ligninger er at $T_c \cdot V_c = 0$.
Så en av disse må være
null.

$\Rightarrow$ Det finnes ingen ikke-trivielle
kritiske punkt

$\Rightarrow$ Nei

$$36) \quad \frac{dP}{dT} = \frac{L_f}{T \Delta V}$$

$$\Delta V = V_g - V_v \approx V_g = \frac{NkT}{P}$$

$$\frac{dP}{dT} = \frac{L_f P}{NkT^2}$$

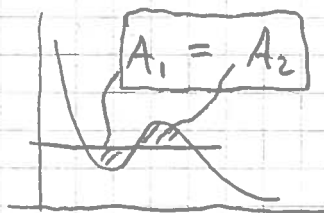
$$\frac{Nk}{L_f} \int_{P_1}^P \frac{dP}{P} = \int_{T_1}^T \frac{dT}{T^2}$$

$$\frac{Nk}{L_f} \ln \frac{P}{P_1} = - \left[\frac{1}{T} \right]_{T_1}^T = - \left[\frac{1}{T} - \frac{1}{T_1} \right]$$

$$\underline{\underline{\ln \frac{P}{P_1} = - \frac{L_f}{Nk} \left[\frac{1}{T} - \frac{1}{T_1} \right]}}$$

37) C

38) Likvektstrykket er bestemt av Maxwell-struksjonen.



$$\underline{\underline{P/P_c = 0.8}}$$

39) $LV \Leftrightarrow \mu_{\text{glass}} = \mu_{\text{kristall}}$

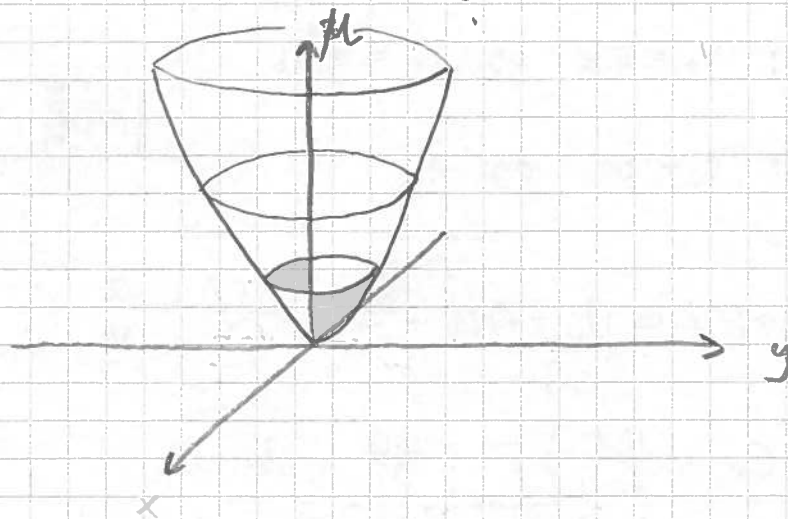
$$\mu_g = \frac{\beta}{2} T^2 - \beta T^2 = -\frac{\beta}{2} T^2$$

$$\mu_k = \frac{\alpha}{4} T^4 - \frac{\alpha}{3} T^4 = -\frac{\alpha}{12} T^4$$

$$\Rightarrow \frac{\beta}{2} T_m^2 = \frac{\alpha}{12} T_m^4$$

$$T_m^2 = \frac{6\beta}{\alpha} \Rightarrow \underline{\underline{T_m = \sqrt{\frac{6\beta}{\alpha}}}}$$

40) Partikler strømmer fra høy til lav μ .



Partiklene strømmer mot origo.

$$41) \quad G = U + pV - TS$$

$$\left(\frac{\partial S}{\partial V}\right)_T = \left(\frac{\partial p}{\partial T}\right)_V = \frac{\partial}{\partial T} \left\{ \frac{a}{3} T^4 \right\} = \frac{4aT^3}{3}$$

$$S(V, T) = \frac{4aT^3}{3} V + \underbrace{C(T)}_{=0}$$

$$G = aVT^4 + \frac{a}{3} VT^4 - \frac{4a}{3} VT^4 = 0!$$

$$\text{Siden } \mu = G/N \Rightarrow \underline{\underline{\mu = 0}}$$

$$42) \quad \text{Før: } V_0 = 5b \quad \text{og} \quad T_0 = \frac{a}{2Rb}$$

$$\text{Slutt: } V_s = \infty \quad \text{og} \quad T_s$$



$$\underline{H} = U + pV = U_0 + C_v T - \frac{a}{V} + \frac{RTV}{V-b} - \frac{a}{V}$$

$$= \left(C_v + \frac{RV}{V-b} \right) T - \frac{2a}{V} + U_0$$

$$\text{Joule-Thomson } \Leftrightarrow H_{\text{før}} = H_{\text{etter}} :$$

$$\left(C_v + R \frac{V_0}{V_0-b} \right) T_0 - \frac{2a}{V_0} = (C_v + R) T_s$$

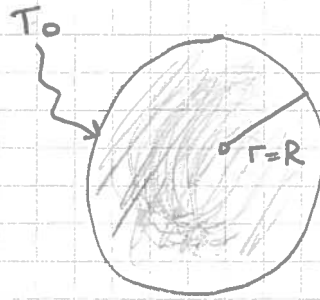
$$\Rightarrow T_s = \frac{39}{50} T_0 \approx \underline{\underline{0.78 T_0}}$$

$$43) \quad \left. \begin{aligned} \vec{J} &= -\chi \nabla T \\ \nabla \cdot \vec{J} &= 0 \quad (\text{stationär}) \end{aligned} \right\} \begin{aligned} J &= -\chi \frac{dT}{dx} \\ \frac{dJ}{dx} &= 0 \end{aligned}$$

$$0 = \frac{d}{dx} \left(-\chi(x) \frac{dT}{dx} \right) = - \left\{ \underbrace{\frac{d\chi}{dx}}_{-e^{-x}} \frac{dT}{dx} + \chi \frac{d^2 T}{dx^2} \right\}$$

$$0 = e^{-x} \frac{dT}{dx} - e^{-x} \frac{d^2 T}{dx^2}$$

$$\Rightarrow \boxed{\frac{d^2 T}{dx^2} - \frac{dT}{dx} = 0}$$



Stationær formeligning med
kulesymmetri:

$$\frac{d}{dr} \left(r^2 \frac{dT}{dr} \right) = 0$$

$$r^2 \frac{dT}{dr} = C_1$$

$$\frac{dT}{dr} = \frac{C_1}{r^2}$$

$$T(r) = C_2 - \frac{C_1}{r}$$

$$T(r \rightarrow \infty) = 0 = C_2$$

$$T(R) = T_0 = -\frac{C_1}{R} \Rightarrow C_1 = -RT_0$$

$$\Rightarrow T(r) = R \left(\frac{T_0}{r} \right)$$

$$45) P(x, t) = \frac{1}{\sqrt{4\pi Dt}} e^{-\frac{(x-v_d t)^2}{4Dt}}$$

$$\langle x \rangle = \int_{-\infty}^{\infty} \frac{1}{\sqrt{4\pi Dt}} e^{-\frac{(x-v_d t)^2}{4Dt}} x dx$$

La : $y = x - v_d t \Rightarrow dy = dx$ og $a_1 = 4Dt$
 $a_2 = 4\pi Dt$

$$\langle x \rangle = \frac{1}{\sqrt{a_2}} \int_{-\infty}^{\infty} (y + v_d t) e^{-\frac{y^2}{a_1}}$$

$$= \underbrace{\frac{1}{\sqrt{a_2}} \int_{-\infty}^{\infty} y e^{-\frac{y^2}{a_1}}}_{0} + \underbrace{\left(\frac{1}{\sqrt{a_2}} \int_{-\infty}^{\infty} e^{-\frac{y^2}{a_1}} \right)}_1 v_d t$$

$$= \underline{\underline{v_d t}} \Rightarrow \text{Partiellere strømmer mot høyre}$$

$$\langle X^2 \rangle = \int_{-\infty}^{\infty} \frac{1}{\sqrt{4\pi Dt}} e^{-\frac{(X-V_{dt})^2}{4Dt}} X^2 dx$$

$$y = X - V_{dt},$$

$$\langle X^2 \rangle = \frac{1}{\sqrt{4\pi Dt}} \int_{-\infty}^{\infty} \underbrace{(y + V_{dt})^2}_{y^2 + 2yV_{dt} + (V_{dt})^2} e^{-\frac{y^2}{4Dt}} dy$$

$$= \frac{1}{\sqrt{4\pi Dt}} \left\{ \int_{-\infty}^{\infty} y^2 e^{-\frac{y^2}{4Dt}} + 2V_{dt} \underbrace{y e^{-\frac{y^2}{4Dt}}}_0 + (V_{dt})^2 \underbrace{e^{-\frac{y^2}{4Dt}}}_1 \right\} dy$$

$$= (V_{dt})^2 + \underbrace{\frac{1}{\sqrt{4\pi Dt}} \int_{-\infty}^{\infty} y^2 e^{-\frac{y^2}{4Dt}} dy}_I$$

$$\lambda = \frac{1}{4Dt}$$

$$I = \int_{-\infty}^{\infty} y^2 e^{-\lambda y^2} dy = \underbrace{-\frac{d}{d\lambda} \int_{-\infty}^{\infty} e^{-\lambda y^2} dy}_{\sqrt{\frac{\pi}{\lambda}}}$$

$$= \frac{1}{2} \sqrt{\pi} \frac{1}{\lambda^{3/2}}$$

$$\langle x^2 \rangle = (V_d t)^2 + \frac{1}{\sqrt{4\pi Dt}} \left(\frac{1}{2} \sqrt{\pi} \sqrt{4Dt} 4Dt \right)$$

$$= \underline{(V_d t)^2 + 2Dt}$$

Vorden
hoger
sey

diffusion

