

Contact during the exam:  
Professor Ingve Simonsen  
Telephone: 9 34 17 or 470 76 416

### Exam in TFY4240 Electromagnetic Theory

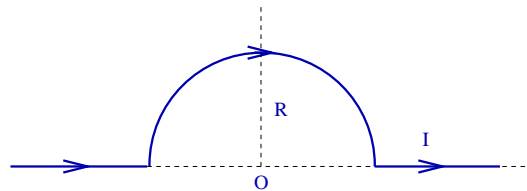
Wednesday Dec 10, 2008  
09:00–13:00

Allowed help: Alternativ C

Authorized calculator and mathematical formula book

This problem set consists of 8=1one page=0.

#### Problem 1.



An infinitely long wire carries a (time-independent) current  $I$ . The wire is bent so as to have a semi-circular detour, of radius  $R$ , around the origin  $O$  (see figure).

- Derive an expression for the magnetic field (vector),  $\mathbf{H}$ , at the origin  $O$  of the coordinate system.
- Determine the numeric value of this magnetic field given the current  $I = 1\text{A}$  and radius  $R = 1\text{cm}$ .

#### Problem 2.

In this problem, we will consider the so-called *attenuation constant* for a plane wave propagating in a good conductor. We aim at, step-by-step, to derive an expression for this constant. The medium under study is an ohmic conductor of permittivity  $\varepsilon$ , permeability  $\mu$  and conductivity  $\sigma$ . For simplicity these constants are assumed to be *independent* of frequency.

- From the Maxwell's equations and Ohm's law, show that the relevant wave equation reads

$$\nabla^2 \mathbf{E} - \mu \varepsilon \partial_t^2 \mathbf{E} - \mu \sigma \partial_t \mathbf{E} = 0, \quad (1)$$

where  $\mathbf{E} \equiv \mathbf{E}(\mathbf{r}, t)$  and  $\partial_t = \frac{\partial}{\partial t}$ .

- b) For a wave of angular frequency  $\omega$ ,  $\mathbf{E}(\mathbf{r}, t) = \mathbf{E}_0(\mathbf{r})e^{-i\omega t}$ , Eq. (1) can be written in the form

$$\nabla^2 \mathbf{E}_0 + \mu\epsilon(\omega)\omega^2 \mathbf{E}_0 = 0. \quad (2)$$

Show this, and identify the function  $\epsilon(\omega)$  (different from  $\epsilon$ ).

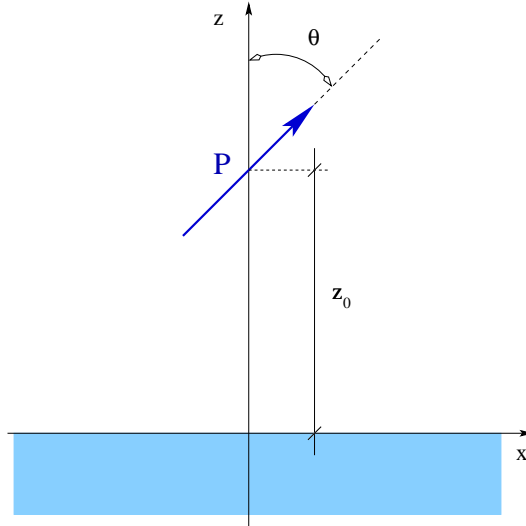
A plane wave is incident on the conductor along the inward normal, whose direction is taken to be the  $z$ -direction. Then in the conductor the electromagnetic wave can be represented by

$$\mathbf{E} = \mathbf{E}_0 e^{ikz - i\omega t}, \quad (3)$$

where  $k$  is the wave number.

- c) Find an expression for the wave number  $k$  in terms of  $\omega$  and the medium parameters ( $\epsilon$ ,  $\mu$  and  $\sigma$ ).
- d) We write  $k = k_1 + ik_2$ , where  $k_1$  and  $k_2$  both are real functions. For a good conductor, *i.e.* for  $\sigma/(\epsilon\omega) \gg 1$ , show that  $k_1 = k_2$  and determine this common function (again) in terms of  $\omega$  and the material parameters.
- e) Argue why it is reasonable to name the constant  $\delta = 1/k_2$  the *attenuation constant*. Write down the expression for this constant ( $\delta$ ).

### Problem 3.



A static electric dipole is located in vacuum at position  $\mathbf{r}_0 = (0, 0, z_0)$  (see figure). Its dipole moment can be written  $\mathbf{p} = p(\sin \theta, 0, \cos \theta)$  where  $\theta$  is the angle between  $\mathbf{p}$  and the positive  $z$ -axis. Initially vacuum is filling the whole space (also the region  $z \leq 0$ ).

- a) Show that the scalar potential for an individual dipole (without the conducting half-space present) can be written as

$$V(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{\hat{\mathbf{R}} \cdot \mathbf{p}}{R^2}. \quad (4)$$

What is the meaning of  $\mathbf{R}$  in this equation? In your proof, you may for simplicity set  $z_0 = 0$  and  $\theta = 0$ . Below, however, this assumption will *not* be made.

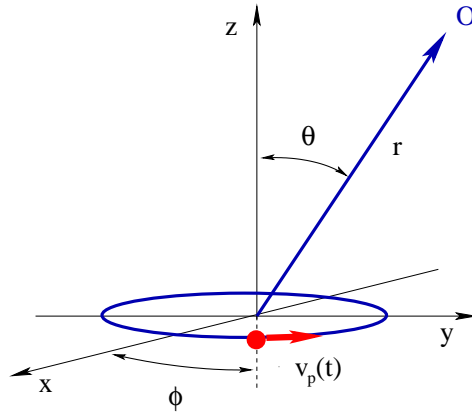
Now a perfectly conducting, *grounded*, half-space is placed at  $z \leq 0$ .

- b) Give the boundary conditions that the scalar potential,  $V(\mathbf{r})$ , satisfies at the interface of the metallic half-space ( $z = 0$ ). Explain (in words) the essence of the method of images.
- c) Use the results from point b) to determine the location and orientation of the image dipole, and make a sketch of the resulting configuration. Moreover, show that the scalar potential for  $z \geq 0$  can be written as

$$V(\mathbf{r}) = \frac{p}{4\pi\epsilon_0} \left( \frac{x \sin \theta + (z - z_0) \cos \theta}{[x^2 + y^2 + (z - z_0)^2]^{3/2}} + \frac{-x \sin \theta + (z + z_0) \cos \theta}{[x^2 + y^2 + (z + z_0)^2]^{3/2}} \right). \quad (5)$$

- d) Determine the induced (surface) charge density,  $\sigma(x, y)$  on the surface of the metal. Express your answer in terms of the spatial coordinates  $x$  and  $y$ , the dipole height  $z_0$ , the dipole orientation  $\theta$ , and the magnitude of the dipole moment  $|\mathbf{p}| = p$ .

#### Problem 4.



Consider a particle of charge  $q \neq 0$  that is moving with a constant angular frequency  $\omega$  along a circular path of radius  $r_0$  in the  $xy$ -plane (see figure). For instance, this can be achieved by applying a static magnetic field  $\mathbf{H}$ . It is assumed that the particle velocity is *non-relativistic* ( $v_p \ll c$ ).

An observer point,  $O$ , is defined by the spherical coordinates  $(r, \theta, \phi)$  relative to a coordinates system with origin in the centered of the circle (see figure).

- a) Write down an expression for the *time-dependent* particle position,  $\mathbf{r}_p(t)$ , and use this to calculate the particle velocity,  $\mathbf{v}_p(t)$ , and acceleration,  $\mathbf{a}_p(t)$ . What is the direction of the applied (static) magnetic field,  $\mathbf{H}$ , relative  $\mathbf{v}_p(t)$ , for the particle to make circular motion?

We will now study the radiation from this particle. The time-dependent radiated power per solid angle is given by

$$\frac{dP}{d\Omega} = \frac{1}{4\pi\epsilon_0} \frac{q^2}{4\pi c^3} \left( \hat{\mathbf{R}} \times \mathbf{a}_p \right)^2. \quad (6)$$

- b) Explain what the time-dependent factor  $\hat{\mathbf{R}}(t)$  in Eq. (6) means. When calculating  $dP/d\Omega$ , what time should be used for this quantity and  $\mathbf{a}_p(t)$ ?
- c) Under the assumption  $r_0 \ll r$  derive an expression for the *time-averaged* radiated power per solid angle  $\langle dP/d\Omega \rangle$  for the particle. Why is this expression independent of the angle  $\phi$ ? Explain why the assumption  $r_0 \ll r$  simplifies significantly the calculation.
- d) What is the total radiated power,  $P$ , from the system (independent of radiation direction)? Compare your result with Larmor's formula (*cf.* formula sheet).
- e) You have just showed (hopefully) that the particle is radiating, *i.e.* that  $P \neq 0$ . However, still the particle performs circular motion of constant angular velocity, and therefore has time independent total energy. Explain how this is possible. Where is the radiated energy coming from?

## FUNDAMENTAL CONSTANTS

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$$\epsilon_0 = 8.85 \times 10^{-12} \text{ C}^2/\text{Nm}^2 \quad (\text{permittivity of free space})$$

$$\mu_0 = 4\pi \times 10^{-7} \text{ N/A}^2 \quad (\text{permeability of free space})$$

$$c = 3.00 \times 10^8 \text{ m/s} \quad (\text{speed of light})$$

$$e = 1.60 \times 10^{-19} \text{ C} \quad (\text{charge of the electron})$$

$$m = 9.11 \times 10^{-31} \text{ kg} \quad (\text{mass of the electron})$$

## SPHERICAL AND CYLINDRICAL COORDINATES

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### Spherical

$$\begin{cases} x = r \sin \theta \cos \phi \\ y = r \sin \theta \sin \phi \\ z = r \cos \theta \end{cases} \quad \begin{cases} \hat{\mathbf{x}} = \sin \theta \cos \phi \hat{\mathbf{r}} + \cos \theta \cos \phi \hat{\theta} - \sin \phi \hat{\phi} \\ \hat{\mathbf{y}} = \sin \theta \sin \phi \hat{\mathbf{r}} + \cos \theta \sin \phi \hat{\theta} + \cos \phi \hat{\phi} \\ \hat{\mathbf{z}} = \cos \theta \hat{\mathbf{r}} - \sin \theta \hat{\theta} \end{cases}$$

$$\begin{cases} r = \sqrt{x^2 + y^2 + z^2} \\ \theta = \tan^{-1}(\sqrt{x^2 + y^2}/z) \\ \phi = \tan^{-1}(y/x) \end{cases} \quad \begin{cases} \hat{\mathbf{r}} = \sin \theta \cos \phi \hat{\mathbf{x}} + \sin \theta \sin \phi \hat{\mathbf{y}} + \cos \theta \hat{\mathbf{z}} \\ \hat{\theta} = \cos \theta \cos \phi \hat{\mathbf{x}} + \cos \theta \sin \phi \hat{\mathbf{y}} - \sin \theta \hat{\mathbf{z}} \\ \hat{\phi} = -\sin \phi \hat{\mathbf{x}} + \cos \phi \hat{\mathbf{y}} \end{cases}$$

### Cylindrical

$$\begin{cases} x = s \cos \phi \\ y = s \sin \phi \\ z = z \end{cases} \quad \begin{cases} \hat{\mathbf{x}} = \cos \phi \hat{\mathbf{s}} - \sin \phi \hat{\phi} \\ \hat{\mathbf{y}} = \sin \phi \hat{\mathbf{s}} + \cos \phi \hat{\phi} \\ \hat{\mathbf{z}} = \hat{\mathbf{z}} \end{cases}$$

$$\begin{cases} s = \sqrt{x^2 + y^2} \\ \phi = \tan^{-1}(y/x) \\ z = z \end{cases} \quad \begin{cases} \hat{\mathbf{s}} = \cos \phi \hat{\mathbf{x}} + \sin \phi \hat{\mathbf{y}} \\ \hat{\phi} = -\sin \phi \hat{\mathbf{x}} + \cos \phi \hat{\mathbf{y}} \\ \hat{\mathbf{z}} = \hat{\mathbf{z}} \end{cases}$$

# BASIC EQUATIONS OF ELECTRODYNAMICS

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## Maxwell's Equations

*In general :*

$$\left\{ \begin{array}{l} \nabla \cdot \mathbf{E} = \frac{1}{\epsilon_0} \rho \\ \nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \\ \nabla \cdot \mathbf{B} = 0 \\ \nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \end{array} \right.$$

*In matter :*

$$\left\{ \begin{array}{l} \nabla \cdot \mathbf{D} = \rho_f \\ \nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \\ \nabla \cdot \mathbf{B} = 0 \\ \nabla \times \mathbf{H} = \mathbf{J}_f + \frac{\partial \mathbf{D}}{\partial t} \end{array} \right.$$

## Auxiliary Fields

*Definitions :*

$$\left\{ \begin{array}{l} \mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P} \\ \mathbf{H} = \frac{1}{\mu_0} \mathbf{B} - \mathbf{M} \end{array} \right.$$

*Linear media :*

$$\left\{ \begin{array}{l} \mathbf{P} = \epsilon_0 \chi_e \mathbf{E}, \quad \mathbf{D} = \epsilon \mathbf{E} \\ \mathbf{M} = \chi_m \mathbf{H}, \quad \mathbf{H} = \frac{1}{\mu} \mathbf{B} \end{array} \right.$$

## Potentials

$$\mathbf{E} = -\nabla V - \frac{\partial \mathbf{A}}{\partial t}, \quad \mathbf{B} = \nabla \times \mathbf{A}$$

## Lorentz force law

$$\mathbf{F} = q(\mathbf{E} + \mathbf{v} \times \mathbf{B})$$

## Energy, Momentum, and Power

$$\text{Energy :} \quad U = \frac{1}{2} \int \left( \epsilon_0 E^2 + \frac{1}{\mu_0} B^2 \right) d\tau$$

$$\text{Momentum :} \quad \mathbf{P} = \epsilon_0 \int (\mathbf{E} \times \mathbf{B}) d\tau$$

$$\text{Poynting vector :} \quad \mathbf{S} = \frac{1}{\mu_0} (\mathbf{E} \times \mathbf{B})$$

$$\text{Larmor formula :} \quad P = \frac{\mu_0}{6\pi c} q^2 a^2$$

## VECTOR IDENTITIES

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### Triple Products

$$(1) \quad \mathbf{A} \cdot (\mathbf{B} \times \mathbf{C}) = \mathbf{B} \cdot (\mathbf{C} \times \mathbf{A}) = \mathbf{C} \cdot (\mathbf{A} \times \mathbf{B})$$

$$(2) \quad \mathbf{A} \times (\mathbf{B} \times \mathbf{C}) = \mathbf{B}(\mathbf{A} \cdot \mathbf{C}) - \mathbf{C}(\mathbf{A} \cdot \mathbf{B})$$

### Product Rules

$$(3) \quad \nabla(fg) = f(\nabla g) + g(\nabla f)$$

$$(4) \quad \nabla(\mathbf{A} \cdot \mathbf{B}) = \mathbf{A} \times (\nabla \times \mathbf{B}) + \mathbf{B} \times (\nabla \times \mathbf{A}) + (\mathbf{A} \cdot \nabla)\mathbf{B} + (\mathbf{B} \cdot \nabla)\mathbf{A}$$

$$(5) \quad \nabla \cdot (f\mathbf{A}) = f(\nabla \cdot \mathbf{A}) + \mathbf{A} \cdot (\nabla f)$$

$$(6) \quad \nabla \cdot (\mathbf{A} \times \mathbf{B}) = \mathbf{B} \cdot (\nabla \times \mathbf{A}) - \mathbf{A} \cdot (\nabla \times \mathbf{B})$$

$$(7) \quad \nabla \times (f\mathbf{A}) = f(\nabla \times \mathbf{A}) - \mathbf{A} \times (\nabla f)$$

$$(8) \quad \nabla \times (\mathbf{A} \times \mathbf{B}) = (\mathbf{B} \cdot \nabla)\mathbf{A} - (\mathbf{A} \cdot \nabla)\mathbf{B} + \mathbf{A}(\nabla \cdot \mathbf{B}) - \mathbf{B}(\nabla \cdot \mathbf{A})$$

### Second Derivatives

$$(9) \quad \nabla \cdot (\nabla \times \mathbf{A}) = 0$$

$$(10) \quad \nabla \times (\nabla f) = 0$$

$$(11) \quad \nabla \times (\nabla \times \mathbf{A}) = \nabla(\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A}$$

## FUNDAMENTAL THEOREMS

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$$\text{Gradient Theorem : } \int_a^b (\nabla f) \cdot d\mathbf{l} = f(\mathbf{b}) - f(\mathbf{a})$$

$$\text{Divergence Theorem : } \int (\nabla \cdot \mathbf{A}) d\tau = \oint \mathbf{A} \cdot d\mathbf{a}$$

$$\text{Curl Theorem : } \int (\nabla \times \mathbf{A}) \cdot d\mathbf{a} = \oint \mathbf{A} \cdot d\mathbf{l}$$

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## VECTOR DERIVATIVES

**Cartesian.**  $d\mathbf{l} = dx \hat{\mathbf{x}} + dy \hat{\mathbf{y}} + dz \hat{\mathbf{z}}; \quad d\tau = dx dy dz$

$$\text{Gradient :} \quad \nabla t = \frac{\partial t}{\partial x} \hat{\mathbf{x}} + \frac{\partial t}{\partial y} \hat{\mathbf{y}} + \frac{\partial t}{\partial z} \hat{\mathbf{z}}$$

$$\text{Divergence :} \quad \nabla \cdot \mathbf{v} = \frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} + \frac{\partial v_z}{\partial z}$$

$$\text{Curl :} \quad \nabla \times \mathbf{v} = \left( \frac{\partial v_z}{\partial y} - \frac{\partial v_y}{\partial z} \right) \hat{\mathbf{x}} + \left( \frac{\partial v_x}{\partial z} - \frac{\partial v_z}{\partial x} \right) \hat{\mathbf{y}} + \left( \frac{\partial v_y}{\partial x} - \frac{\partial v_x}{\partial y} \right) \hat{\mathbf{z}}$$

$$\text{Laplacian :} \quad \nabla^2 t = \frac{\partial^2 t}{\partial x^2} + \frac{\partial^2 t}{\partial y^2} + \frac{\partial^2 t}{\partial z^2}$$

**Spherical.**  $d\mathbf{l} = dr \hat{\mathbf{r}} + r d\theta \hat{\boldsymbol{\theta}} + r \sin \theta d\phi \hat{\boldsymbol{\phi}}; \quad d\tau = r^2 \sin \theta dr d\theta d\phi$

$$\text{Gradient :} \quad \nabla t = \frac{\partial t}{\partial r} \hat{\mathbf{r}} + \frac{1}{r} \frac{\partial t}{\partial \theta} \hat{\boldsymbol{\theta}} + \frac{1}{r \sin \theta} \frac{\partial t}{\partial \phi} \hat{\boldsymbol{\phi}}$$

$$\text{Divergence :} \quad \nabla \cdot \mathbf{v} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 v_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta v_\theta) + \frac{1}{r \sin \theta} \frac{\partial v_\phi}{\partial \phi}$$

$$\begin{aligned} \text{Curl :} \quad \nabla \times \mathbf{v} = & \frac{1}{r \sin \theta} \left[ \frac{\partial}{\partial \theta} (\sin \theta v_\phi) - \frac{\partial v_\theta}{\partial \phi} \right] \hat{\mathbf{r}} \\ & + \frac{1}{r} \left[ \frac{1}{\sin \theta} \frac{\partial v_r}{\partial \phi} - \frac{\partial}{\partial r} (r v_\phi) \right] \hat{\boldsymbol{\theta}} + \frac{1}{r} \left[ \frac{\partial}{\partial r} (r v_\theta) - \frac{\partial v_r}{\partial \theta} \right] \hat{\boldsymbol{\phi}} \end{aligned}$$

$$\text{Laplacian :} \quad \nabla^2 t = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial t}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial t}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 t}{\partial \phi^2}$$

**Cylindrical.**  $d\mathbf{l} = ds \hat{\mathbf{s}} + s d\phi \hat{\boldsymbol{\phi}} + dz \hat{\mathbf{z}}; \quad d\tau = s ds d\phi dz$

$$\text{Gradient :} \quad \nabla t = \frac{\partial t}{\partial s} \hat{\mathbf{s}} + \frac{1}{s} \frac{\partial t}{\partial \phi} \hat{\boldsymbol{\phi}} + \frac{\partial t}{\partial z} \hat{\mathbf{z}}$$

$$\text{Divergence :} \quad \nabla \cdot \mathbf{v} = \frac{1}{s} \frac{\partial}{\partial s} (s v_s) + \frac{1}{s} \frac{\partial v_\phi}{\partial \phi} + \frac{\partial v_z}{\partial z}$$

$$\text{Curl :} \quad \nabla \times \mathbf{v} = \left[ \frac{1}{s} \frac{\partial v_z}{\partial \phi} - \frac{\partial v_\phi}{\partial z} \right] \hat{\mathbf{s}} + \left[ \frac{\partial v_s}{\partial z} - \frac{\partial v_z}{\partial s} \right] \hat{\boldsymbol{\phi}} + \frac{1}{s} \left[ \frac{\partial}{\partial s} (s v_\phi) - \frac{\partial v_s}{\partial \phi} \right] \hat{\mathbf{z}}$$

$$\text{Laplacian :} \quad \nabla^2 t = \frac{1}{s} \frac{\partial}{\partial s} \left( s \frac{\partial t}{\partial s} \right) + \frac{1}{s^2} \frac{\partial^2 t}{\partial \phi^2} + \frac{\partial^2 t}{\partial z^2}$$