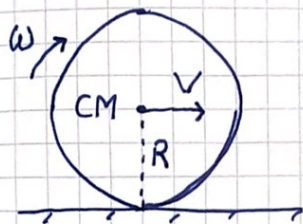


Ren rulling

[051 11.1]

(24)



For hel omdreining: $V = \frac{2\pi R}{T}$; $\omega = \frac{2\pi}{T}$

$$\Rightarrow \underline{V = \omega R}$$

Rullebetingelsen

$$K = \frac{1}{2}MV^2 + \frac{1}{2}I_0\omega^2 = \frac{1}{2}MV^2 + \frac{1}{2} \cdot cMR^2 \cdot \frac{V^2}{R^2} = \underline{\frac{1}{2}(1+c)MV^2}$$

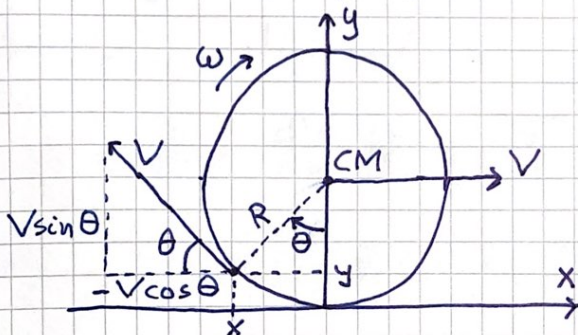
Farten ute på periferien:

$$\omega = \dot{\theta} = \frac{V}{R} \Rightarrow V = R\dot{\theta}$$

Fra figuren:

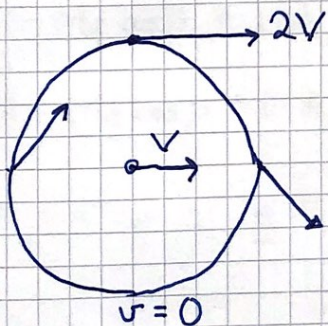
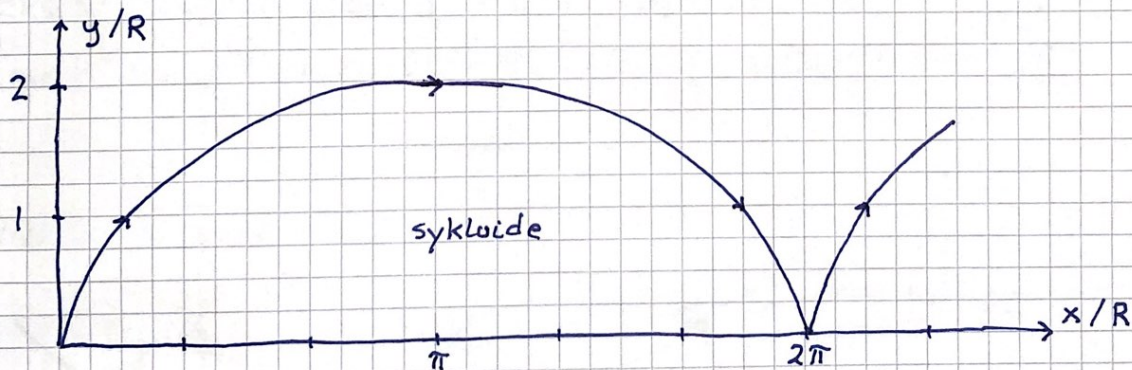
$$v_x = V - V\cos\theta$$

$$v_y = V\sin\theta$$



Barer til punkt på periferien; velger $x=y=0$ for $\theta=0$:

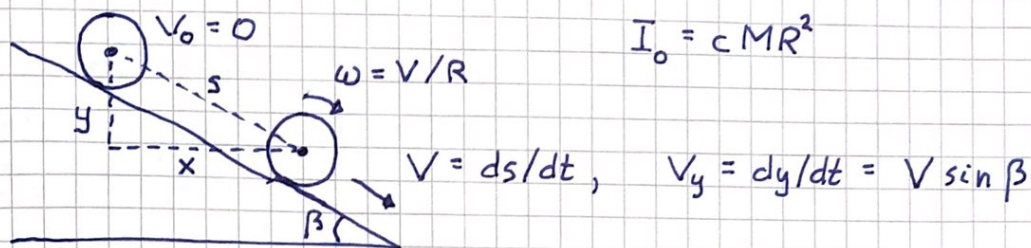
$$x(\theta) = R\theta - R\sin\theta \quad ; \quad y(\theta) = R - R\cos\theta$$



Kontaktpunktet er i ro,
hjulet glir ikke.

Eks: Ren rulling på skråplan

(25)



Bestem V , A , f og $\max \beta$ som gir ren rulling

Løsning: Statisk $f \Rightarrow$ Mek. energi bevart

$$\Rightarrow Mgy = \frac{1+c}{2} MV^2 \Rightarrow \underline{V(y) = \sqrt{\frac{2gy}{1+c}}}$$

$$A = \frac{dV}{dt} = \frac{dV}{dy} \frac{dy}{dt} = \sqrt{\frac{2g}{1+c}} \cdot \frac{1}{2} y^{-1/2} \cdot \sqrt{\frac{2gy}{1+c}} \sin \beta = \underline{\frac{g \sin \beta}{1+c}}$$

$$N2: Mg \sin \beta - f = MA \Rightarrow \underline{f = \frac{c}{1+c} Mg \sin \beta}$$

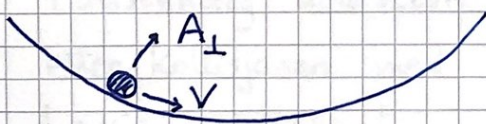
$$f \leq \mu_s N ; \quad \underline{N = Mg \cos \beta}$$

$$\Rightarrow \frac{c}{1+c} Mg \sin \beta \leq \mu_s Mg \cos \beta \Rightarrow \underline{\tan \beta \leq \mu_s \cdot \frac{1+c}{c}}$$

Eks: Kompakt kule ($c=2/5$); $\mu_s = 1/3$. Hva er $\beta_{\max}$?

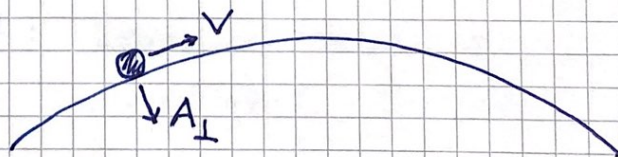
$$\text{Løsn: } \beta_{\max} = \arctan\left\{\frac{(1+2/5)}{(3 \cdot 2/5)}\right\} = \arctan(7/6) = \underline{49^\circ}$$

Rulling på krum bane:



$$N - Mg \cos \beta = MA_{\perp}$$

$$\Rightarrow \underline{N = Mg \cos \beta + MA_{\perp}}$$



$$Mg \cos \beta - N = MA_{\perp}$$

$$\Rightarrow \underline{N = Mg \cos \beta - MA_{\perp}}$$

$$\underline{A_{\perp} = V^2/\rho} ; \quad \rho = \text{krumningsradien (se s. 6)}$$

Impuls, Kollisjoner, Rakettprinsipp [os1 9]

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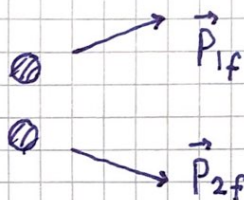
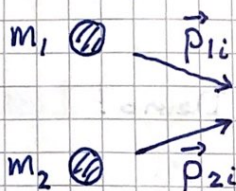
N2 for masse m : $\vec{F} = m \frac{d\vec{v}}{dt} = \frac{d}{dt}(m\vec{v})$

Massens impuls: $\vec{p} = m\vec{v}$; $[p] = \text{kg} \cdot \text{m/s}$

$\Rightarrow \vec{F} = d\vec{p}/dt$

Impulsbevarelse: Hvis $\vec{F} = 0$ er $\vec{p}$ bevart

Kollisjoner:



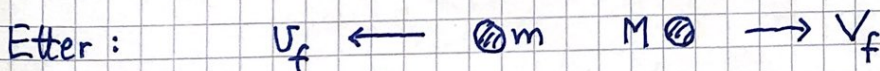
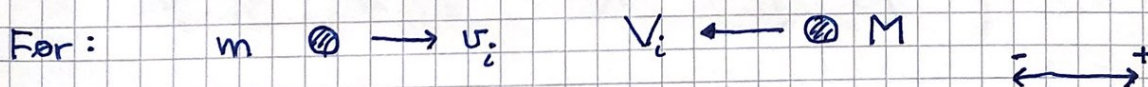
N3 $\Rightarrow \vec{F}_{21} + \vec{F}_{12} = 0 \xrightarrow{\text{N2}} d\vec{p}_1/dt + d\vec{p}_2/dt = 0$

$\Rightarrow \vec{p}_1 + \vec{p}_2$ er konstant dersom ytre krefter kan neglisjeres i kollisjonen (typisk kortvarig)

Kinetisk energi ($K = \sum_i K_i$) kan gå tapt pga friksjon og deformasjon.

- Elastisk støt : $\Delta K = 0$
- Uelastisk støt : $\Delta K < 0$
- Fullstendig uelastisk støt: Legemene henger sammen etter kollisjonen med felles slutthastighet. Da er $|\Delta K|$ maksimal.

Vi ser kun på et sentralt støt, dvs kollisjon i 1D:



$$\Delta p = 0 : \quad m u_i + M V_i = m u_f + M V_f \quad (1) \quad (27)$$

Hvis fullstendig uelastisk : $u_f = V_f = \frac{m u_i + M V_i}{m + M}$

Elastisk : $\frac{1}{2} m u_i^2 + \frac{1}{2} M V_i^2 = \frac{1}{2} m u_f^2 + \frac{1}{2} M V_f^2 \quad (2)$

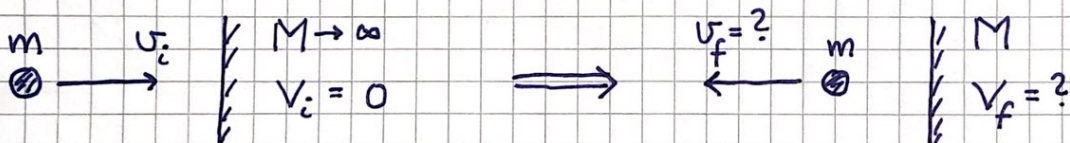
Løsning av (1) og (2) mhp slutt hastighetene gir:

$$u_f = \frac{M}{m+M} \left(2V_i + u_i \cdot \frac{m-M}{M} \right) ; \quad V_f = \frac{m}{M+m} \left(2u_i + V_i \cdot \frac{M-m}{m} \right)$$

Eks 1: $M = m, V_i = 0$ (Newtons rulle)

Da er $u_f = 0$ og $V_f = u_i$

Eks 2: Ball mot vegg



Løsning 2:

$$u_f = \frac{M}{m+M} \cdot u_i \cdot \frac{m-M}{M} = \underline{\underline{-u_i}} ; \quad V_f = \frac{2m}{M+m} \cdot u_i = \underline{\underline{0}}$$

Impulsbevarelse:

$$\underline{m u_i} = m u_f + M V_f = -m u_i + \frac{2m u_i}{M+m} \cdot M = -m u_i + 2m u_i = \underline{m u_i}$$

Kraftstøt (eng: impulse):

$$\vec{F} = d\vec{p}/dt \Rightarrow d\vec{p} = \vec{F} \cdot dt \Rightarrow \Delta\vec{p} = \int d\vec{p} = \int \vec{F}(t) dt$$